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# A theoretical analysis of zero-field splitting parameters of Mn<sup>2+</sup> doped dicadmium diammonium sulfate Cd<sub>2</sub>(NH<sub>4</sub>)<sub>2</sub>(SO<sub>4</sub>)<sub>3</sub> single crystal

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#### ABSTRACT

The superposition model (SPM) and crystallographic data are utilized to determine the zero-field splitting (ZFS) parameters (ZFSPs) for  $Mn^{2+}$  ions in  $Cd_2(NH_4)_2(SO_4)_3$  single crystal, assuming that the  $Mn^{2+}$  ions locate at either  $Cd^{2+}$  or  $NH^{4+}$  site. The SPM results has been verified by the fourth-order perturbation formulae analysis. Experimental suggestions about  $Mn^{2+}$  ions substituting at  $Cd^{2+}$  sites have been confirmed theoretically for the first time.

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### 1. Introduction

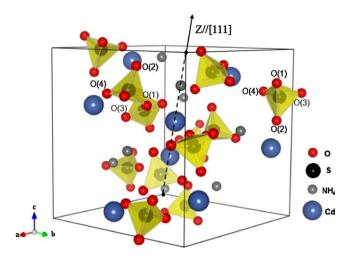
Dicadmium diammonium sulfate Cd<sub>2</sub>(NH<sub>4</sub>)<sub>2</sub>(SO<sub>4</sub>)<sub>3</sub> crystal, abbreviated as CAS hereafter, is one of the langbeinite-type crystals having the general formula  $X_2Y_2(SO_4)_3$ , where X is a divalent metal, e.g. Cd, Zn, Mg, Ca and Y is ammonium (NH<sub>4</sub>) or a monovalent metal, e.g. K, Tl, Rb, Cs [1,2]. CAS has been the subject of many investigations [3–10] since its discovery by Jona and Pepinsky [3]. The ferroelectricity of CAS below 90 K was found out as well [3]. The disordering of the SO<sub>4</sub> ion in the high temperature phase was suggested based on Raman, infrared and far-infrared studies since no mode softening was observed [4]. Electron paramagnetic resonance (EPR) investigations of Mn<sup>2+</sup> doped CAS single-crystal were reported in [5–7]; first at room temperature [5], whereas more detailed study [6] included also the crystal structure of Mn<sup>2+</sup>:CAS. No change in space group symmetry at 80 K could be observed [6], in spite of the fact that a monoclinic ferroelectric phase of CAS below 90 K was suggested [2,3]. The structural phase transitions in CAS were studied by EPR in [7] and [8,9]. The studies [7–10], reporting a proton NMR, indicated some dynamical disordered arrangements related to ammonium and sulfate ions. Also, the authors [7] suggested that the phase transition occurred because of the freezing of the sulfate rotations, while the structural change was considered as a result of some motional change of the sulfate group in [6]. The zero-field splitting (ZFS) parameters (ZFSPs) of  $\rm Mn^{2+}$  in CAS were measured by EPR at room temperature [5–7].

In the present paper, first theoretical analysis of the ZFSPs of  $\mathrm{Mn^{2^+}}$  ions in CAS has been carried out. The ZFSPs have been calculated for different paramagnetic centers formed by  $\mathrm{Mn^{2^+}}$  ions at possible sites, namely,  $\mathrm{Cd^{2^+}}$  site and  $\mathrm{NH_4^+}$  site, in CAS crystal. We employ two theoretical methods: the superposition model (SPM) [11–13] and the fourth-order perturbation formulae on the basis of the dominant spin–orbit coupling mechanism [14]. The outcomes of both methods are consistent with the experimental results and they support earlier suggestions [5–7] that  $\mathrm{Mn^{2^+}}$  ions substitute for  $\mathrm{Cd^{2^+}}$  ions in CAS crystal.

#### 2. Crystal structure

CAS crystal has cubic symmetry with the space group  $P2_13$  ( $T^4$ ) and contain four formula units per unit cell at room temperature. The cubic lattice parameters were reported to be a=1.0350 nm [3,2] and 1.0362 nm [6]. Positional parameters and thermal parameters of CAS at room temperature were reported in [2,6]. Single unit cell of CAS is shown in Fig. 1 with the orientation of the crystallographic axes a, b, and c. The body diagonal of the cubic unit cell is parallel to the symmetry Z-axis [5] (see Fig. 1). CAS undergoes a first-order phase transition from the high-temperature cubic phase ( $P2_1$ ) to the low-temperature monoclinic phase ( $P2_1$ ) [2,15]. CAS is somorphous to the langbeinite  $K_2Mg_2(SO_4)_3$ . The  $NH_4^+$  and  $Cd^{2+}$  lie on the threefold axes, and the  $SO_4$  tetrahedra are in general positions

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**Fig. 1.** Unit cell of  $Cd_2(NH_4)_2(SO_4)$  (CAS) crystal. Orientation of the crystallographic axes a, b, and c and the symmetry axis (Z-axis) parallel to the body diagonal of the cubic unit cell are shown.

**Table 1**Bond distances and coordination angles of ligands used for SPM calculations of ZFS parameters for  $Mn^{2+}$  ions at  $Cd^{2+}$  site in CAS.

Host Cd sites		Substitutional si	Substitutional sites		
Metal-oxygen	R <sub>hi</sub> (nm)	Dopant Mn <sup>2+</sup> -oxygen	$R_i$ (nm)	$\theta_i$ (°)	
Cd-O(1)	0.382873	Mn-O(1)	0.368873	46.7902	
Cd-O(2)	0.425824	Mn-O(2)	0.411824	36.0849	
Cd-O(3)	0.451958	Mn-O(3)	0.437958	64.8707	
Cd-O(4)	0.224573	Mn-O(4)	0.210573	59.8448	

**Table 2** Bond distances and coordination angles of ligands used for SPM calculations of ZFS parameters for  $Mn^{2+}$  ions at  $NH_4^+$  site in CAS.

Host NH <sub>4</sub> sites		Substitutional si	Substitutional sites		
Metal-oxygen	R <sub>hi</sub> (nm)	Dopant Mn <sup>2+</sup> -oxygen	$R_i$ (nm)	$\theta_i$ (°)	
NH <sub>4</sub> -O(1) NH <sub>4</sub> -O(2) NH <sub>4</sub> -O(3) NH <sub>4</sub> -O(4)	0.367858 0.29626 0.513286 0.434782	Mn-O(1) Mn-O(2) Mn-O(3) Mn-O(4)	0.327358 0.25576 0.472786 0.394282	49.3410 57.8399 52.8615 26.5269	

#### 3. Theoretical analysis

#### 3.1. The coordinates of the ligand ions

The bond distances and corresponding coordination angles of ligands have been determined and tabulated for host Cd site in Table 1 and for host NH<sub>4</sub> site in Table 2. The distance of a ligand is usually different from the cation–anion distance in the host lattice because of the mismatch of the radius of the substitution atom  $r_s$  and that of the host atom  $r_h$ . Ligand distance can be reasonably approximated using the following formula [16]:  $R_i \approx R_{hi} + (1/2)(r_s - r_h)$ , where  $R_i$  and  $R_{hi}$  represent the ligand and cation–anion distances, respectively. Using the data:  $r_i(\text{Mn}^{2+}) = 0.067$ ,  $r_h(\text{Cd}^{2+}) = 0.092$  [17], and  $r_h(\text{NH}_4^+) = 0.148$  [18] (in nm), and the cation–anion distances in Tables 1 and 2 yield  $R_i$  values given in Tables 1 and 2.

#### 3.2. Superposition model analysis

Experimental spectra of  $Mn^{2+}$  doped CAS can be analyzed by utilizing the spin Hamiltonian, suitable for the spin S = 5/2 systems at trigonal type I symmetry sites (D<sub>3</sub>, C<sub>3v</sub>, D<sub>3d</sub>), consisting of the

Zeeman electronic terms and the ZFS terms (without considering the hyperfine terms) [19–22]:

$$H = \mu_B \mathbf{B} \cdot \mathbf{g} \cdot \mathbf{S} + \sum f_k b_k^q O_k^q = g_\perp \mu_B (B_x S_x + B_y S_y) + g_\parallel \mu_B B_z S_z$$
$$+ f_k (b_2^0 O_2^0 + b_4^0 O_4^0 + b_4^3 O_4^3) \tag{1}$$

where  $\mathbf{g}$  is the spectroscopic splitting factor,  $\mu_B$  – Bohr magneton,  $\mathbf{B}$  – the applied magnetic field,  $\mathbf{S}$  – the effective spin operator, and  $b_k^q$  are ZFSPs associated with the extended Stevens operators  $O_k^q$ , whereas  $f_k$  = 1/3, and 1/60 are the scaling factors for k = 2, and 4, respectively [23,24].

ZFSPs in Eq. (1) can be estimated using SPM approach outlined recently in [25,26]. The explicit SPM expressions for ZFSPs for a fourfold coordinated  $3d^5$  ion in the Mn-O<sub>4</sub> complex in CAS are derived as

$$\begin{split} D_{SPM} &= b_2^0 = \frac{\bar{b}_2(R_0)}{2} \left[ \left( \frac{R_0}{R_1} \right)^{t_2} (3\cos^2\theta_1 - 1) \right. \\ &+ \left( \frac{R_0}{R_2} \right)^{t_2} (3\cos^2\theta_2 - 1) + \left( \frac{R_0}{R_3} \right)^{t_2} (3\cos^2\theta_3 - 1) \\ &+ \left( \frac{R_0}{R_4} \right)^{t_2} (3\cos^2\theta_4 - 1) \right] \end{split} \tag{2}$$

$$b_4^0 = \frac{\bar{b}_4(R_0)}{8} \left[ \left( \frac{R_0}{R_1} \right)^{t_4} (35\cos^4\theta_1 - 30\cos^2\theta_1 + 3) + \left( \frac{R_0}{R_2} \right)^{t_4} (35\cos^4\theta_2 - 30\cos^2\theta_2 + 3) + \left( \frac{R_0}{R_3} \right)^{t_4} (35\cos^4\theta_3 - 30\cos^2\theta_3 + 3) + \left( \frac{R_0}{R_4} \right)^{t_4} (35\cos^4\theta_4 - 30\cos^2\theta_4 + 3) \right]$$
(3)

$$\begin{split} b_4^3 &= 35 \bar{b}_4(R_0) [\left(\frac{R_0}{R_1}\right)^{t_4} \sin^3\theta_1 \cos\theta_1 + \left(\frac{R_0}{R_2}\right)^{t_4} \sin^3\theta_2 \cos\theta_2 \\ &+ \left(\frac{R_0}{R_3}\right)^{t_4} \sin^3\theta_3 \cos\theta_3 + \left(\frac{R_0}{R_4}\right)^{t_4} \sin^3\theta_4 \cos\theta_4] \end{split}$$

Two data sets [27,28] of the model parameters, i.e. the intrinsic parameters  $\overline{b_k}(R_0)$  and the power law exponents  $t_k$ , suitable for Mn<sup>2+</sup> in CAS exist in literature (Table 5). These sets were adopted in our calculations together with three values of the reference distance  $R_0$  taken as: (a)  $R_0 = a_0/4$ , where  $a_0$  is the lattice parameter, (b)  $R_0$  values from [27,28] for Mn<sup>2+</sup>, and (c)  $R_0 \approx R_{avg}$  [29,30].

#### 3.3. The fourth-order perturbation formula of ZFS parameter D

Second-rank ZFS parameter  $b_2^0 = D$  can also be calculated using the fourth-order perturbation formula on the basis of the dominant spin-orbit coupling mechanism [14]

$$D_{PT} = b_2^0 = \frac{3\zeta^2}{70p^2d}(-B_{20}^2 - 21\zeta B_{20}) + \frac{\zeta^2}{126p^2g}(-10B_{40}^2 + 7B_{43}^2)$$
 (4)

with

$$p = 7B + 7C$$
  $g = 10B + 5C$   $d = 17B + 5C$  (5)

where B and C are the Racah parameters.  $B_{kq}$  are the crystal-field parameters (CFPs) in Wybourne notation. They can be expressed using the superposition model as [14,31]:

$$B_{kq} = \sum_{j} \bar{A}_k(R_j) K_{kq}(\theta_j, \phi_j)$$
 (6)

**Table 3**Electrostatic parameters, spin–orbit coupling coefficient, cubic crystal-field parameter, and orbital reduction factor of Mn<sup>2+</sup> ion at room temperature.

$B_0$	$C_0$	$\zeta_0$	В	С	ζ	Dq	$k \approx N^2$
960 [44,45]	3325 [44,45]	347 [44,45]	690 [43]	3600 [43]	327.8 <sup>a</sup> 300 [6]	1100 <sup>a</sup> 1010 [43] 900 [6]	0.944ª

a This study.

**Table 4**ZFS parameters for Mn<sup>2+</sup>:CAS crystal at room temperature.

Original ZFSPs			Ref.	Converted ZFSPs in units of $[10^{-4}\mathrm{cm}^{-1}]$			
$b_2^0$	$b_4^0$	$b_4^3$		$b_2^0$	$b_4^0$	$b_4^3$	
$4850 \pm 20 \text{GHz}$ $-4719 \pm 4 \text{GHz}$ $-171 \pm 5 \text{G (site 1)}$ $-162 \pm 5 \text{G (site 2)}$	$50 \pm 10 \text{GHz}$ $45 \pm 30 \text{GHz}$ $-1.43 \pm 0.1 \text{G}$ $-1.60 \pm 0.1 \text{G}$	−300 ± 200 GHz 2160 ± 1070 GHz −	[7] [5] [6]	$161.83 \pm 0.67$ $-157.46 \pm 0.13$ $-159.7 \pm 0.5$ $-151.3 \pm 0.5$	$1.67 \pm 0.33$ $1.50 \pm 1$ $-1.34 \pm 0.1$ $-1.49 \pm 0.1$	$-10.01 \pm 6.67$ $72.07 \pm 35.7$	

**Table 5** The adopted SPM parameters and the calculated ZFSPs  $b_{\nu}^{q}$  (in 10<sup>-4</sup> cm<sup>-1</sup>) for the Mn<sup>2+</sup> ions at the Cd<sup>2+</sup> and NH<sub>4</sub>+ sites in CAS.

N / - d - l	(MD)		(1) [27]	C-+ (!!) [27]			
Model parame	eters (MP)	Set (i) [27]		Set (ii) [27]			
$t_2$		8		8			
$\frac{t_4}{b_2}(R_0)$		14		14			
$b_2(R_0)$		-2		$250 \pm 20$			
$\bar{b}_4(R_0)$		0.8		$0.7\pm3$			
$R_0$ (for $\bar{b}_2$ ) (nm			200	0.2101			
$R_0$ (for $\bar{b}_4$ ) (nm	1)	0.2	150	0.2000			
		ZFS parameter	s for Mn <sup>2+</sup> ions at Cd <sup>2+</sup> s	ites			
ZFSPs		$b_2^0$		$b_4^0$		$b_4^3$	
MP sets		i	ii	i	ii	i	ii
Calc.a	*	94.2	-90.5	-1.7	-1.5	66.1	57.8
	**	158.9	<b>-152.8</b>	<b>-4.2</b>	<b>-3.7</b>	162.8	142.4
Calc.b	*	25.7	-17.1	-0.1	-0.04	4.9	1.6
	**	43.4	-28.9	-0.3	-0.1	12.2	3.9
Calc.c	*	1693.0	-1627.9	-267.3	-233.9	10,377.8	9080.6
	**	2857.3	-2747.4	-658.1	<b>-575.9</b>	25,550.0	22,356.2
		ZFS parameter	rs for Mn <sup>2+</sup> ions at NH <sub>4</sub> +	sites			
		$\overline{b_2^0}$		$b_4^0$		$b_4^3$	
MP sets		i	ii	i	ii	i	ii
Calc.a	*	1.6	-1.5	-0.04	-0.04	1.42	1.24
	**	9.6	-9.3	-0.33	-0.29	10.94	9.58
Calc.b	*	0.42	-0.28	-0.003	-0.001	0.11	0.03
	**	2.6	-1.75	-0.02	-0.008	0.82	0.26
Calc.c	*	27.9	-26.9	-6.7	-5.9	222.7	194.9
	**	172.9	-166.3	-51.6	-45.2	1717.9	1503.2

<sup>&</sup>lt;sup>a</sup>  $R_0$  is chosen to be  $a_0/4 \approx 0.25875$  nm.

where  $K_k^q(\theta_i,\phi_i)$  are the coordination factors which are functions of the position angles  $\theta_i$  and  $\phi_i$  of ligands. Here it should be noted that the coordination factors expressed in the Wybourne notation should be distinguished from those in the extended Stevens operator notation [32], which are used to derive the SPM expressions for ZFSPs. For a discussion of the intricate aspects involved in the SPM equations expressed in the two notations, the readers may consult the recent paper [33].  $\bar{A}_k(R_j)$  are the intrinsic parameters and they obey the following power law:  $\bar{A}_k(R_j) = \bar{A}_k(R_0)(R_0/R_j)^{t_k}$ .  $R_0$  and  $R_i$  are the reference distance and the distance of the ith ligand, respectively.  $t_k$  are the power-law exponents.

The explicit SPM expressions for the CFPs  $B_{kq}$  of Mn<sup>2+</sup> ion at the Cd<sup>2+</sup> and NH<sup>4+</sup> sites in CAS have the same mathematical structure with the SPM equations for ZFSPs derived as in the previous section. Intrinsic parameter  $\bar{A}_4$  for CFPs  $B_{40}$  and  $B_{43}$  can be estimated from the cubic CFP Dq from the relation [14]:

$$Dq = -\frac{3}{14}\bar{A}_4[(35\cos^4\theta_{avg} - 30\cos^2\theta_{avg} + 3)$$
$$-7\sqrt{2}\cos\theta_{avg}\sin^3\theta_{avg}] \tag{7}$$

where the averaged values  $\theta_{avg}$  = 51.8977° for Cd<sup>2+</sup> site and  $\theta_{avg}$  = 46.6423° for NH<sub>4</sub><sup>+</sup> site were adopted. Thus, we obtain

<sup>&</sup>lt;sup>b</sup>  $R_0$  as adopted earlier in their source papers (set (i) and set (ii)).

<sup>&</sup>lt;sup>c</sup> Based on the approximation  $R_0 \approx R_{avg} = 0.3713$ .

<sup>\*</sup> Calculations on host crystal structure data.

<sup>\*\*</sup> Calculations based on the contribution from the mismatch of the radii of substitution atom  $r_s$  and host atom  $r_h$ .

 $Dq=1.35\bar{A}_4$  for Mn<sup>2+</sup> at Cd<sup>2+</sup> site, whereas for NH<sub>4</sub><sup>+</sup> site  $Dq=1.28\bar{A}_4$ . The well-known ratio  $\bar{A}_2(R_0)/\bar{A}_4(R_0)\approx 10.8$  for iron group ions in several crystals [34–36] is also used. The following values are adopted:  $t_2=3$ ,  $t_4=5$  as the most common ones in many studies, e.g. [37,38]. Due to the covalence reduction effect for  $3d^n$  ions in crystals both electrostatic parameters B, C and the spin–orbit coupling parameter  $\zeta$  are smaller than the corresponding free-ion values  $(B_0, C_0, \text{and } \zeta_0)$  [39,40]. Orbital reduction factor can be determined from the relation  $N^2 \approx \left(\sqrt{B/B_0} + \sqrt{C/C_0}\right)/2$  and then the spin–orbit coupling parameter is evaluated by  $\zeta \approx N^2 \zeta_0 \approx k$  [41,42]. All relevant parameters adopted for CAS:Mn<sup>2+</sup> system [43] at room temperature are tabulated in Table 3.

#### 4. Results and discussion

The experimental values in the original units and the corresponding ones reconverted to units of  $(10^{-4}) \, \mathrm{cm}^{-1}$  are given in Table 4. Using Eqs. (2) and (3), SPM calculations were performed considering contributions from the four  $\mathrm{O}^{-2}$  ligands around  $\mathrm{Mn}^{2+}$  ions substituting at the  $\mathrm{Cd}^{2+}$  and  $\mathrm{NH}^{4+}$  sites.

The ZFSPs predicted for  $Mn^{2+}$  ions using two modeling approaches, the host crystal structure data (\*) and contributions from the mismatch of the radii of the substitution ( $r_s$ ) and host atom ( $r_h$ ) (\*\*), are listed in Table 5. Comparison of the ZFSPs calculated by SPM (Table 5) with the experimental values (Table 4), indicates that  $Mn^{2+}$  ions locate at the  $Cd^{2+}$  sites. The best values of ZFSPs are obtained when the reference distance  $R_0$  is chosen to be  $a_0/4$ . Generally, taking into account the contribution from the mismatch in radii of the substitution and host atoms yields better results.

There were some inconsistencies for the signs of ZFSPs and the value of  $b_4^3$  in the experimental data of Mn<sup>2+</sup>:CAS (see Table 4). A discrepancy concerning the experimental sign of  $b_2^0(=D)$  exists in the literature; positive in [7] but negative in [5,6]. It was pointed out in [6] that the sign of D and (a-F) was relative, however, the sign of D was supposed negative from a point-charge calculation [46,47]. That is why we utilized two model parameter (MP) sets - they yield similar qualitative results but opposite sign of  $b_2^0$ . The first MP set (i) with negative  $\overline{b_2}$  yields positive  $b_2^0$ , whereas the second MP set (ii) with positive  $\overline{b_2}$  results in negative  $b_2^0$ . For the fourth-rank ZFSPs both MP sets yield the same signs; negative for  $b_4^0$  and positive for  $b_4^3$ . The modeling approach which yields the best result for  $b_2^0$ (bold line in Table 5), results in  $b_{\Delta}^{q}$  differing from the experimental values. However, by just varying  $\bar{b}_4$  to  $0.34 \times 10^{-4}$  cm $^{-1}$  using the flexibility in MP set (ii), we obtain  $b_4^0 = -1.78 \times 10^{-4}$  cm $^{-1}$  and  $b_4^3 = -1.78 \times 10^{-4}$  cm $^{-1}$  cm $^{-1}$  and  $b_4^3 = -1.78 \times 10^{-4}$  cm $^{-1}$  $69.2\times 10^{-4}\,\text{cm}^{-1}\text{,}$  which are in good agreement with the results of [5] (see Table 4). Hence, for Mn<sup>2+</sup> ions in langbeinite-type crystals,  $b_4 = 0.34 \times 10^{-4} \text{ cm}^{-1}$  with  $t_4 = 14$  seem reasonable.

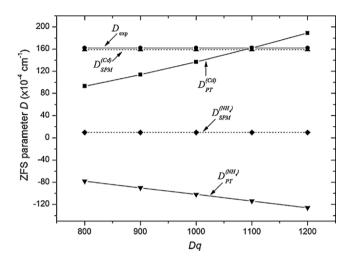
The results of SPM analysis reveal clearly that  $\mathrm{Mn^{2^+}}$  ions can replace  $\mathrm{Cd^{2^+}}$  ions in CAS crystal. In order to confirm this finding and clarify the discrepancy in the sign of  $b_2^0$ , we also derive the secondrank ZFS parameter  $b_2^0(=D)$  using the fourth-order perturbation theory (PT) formula [14]. Hence, the crystal-field (CF) parameters and the zero-field splitting (ZFS) parameter  $D_{PT}$  relevant to the  $\mathrm{Mn^{2^+}}$  ions in tetrahedral coordination at both  $\mathrm{Cd^{2^+}}$  and  $\mathrm{NH_4^+}$  sites in CAS crystal have been calculated (Table 6).

PT calculations yield positive D. Two cubic CF parameter Dq for  $Mn^{2+}$ :CAS exist in literature (see Table 3). This makes us consider Dq as an adjustable parameter. In PT calculations we used  $Dq = 1010 \, \mathrm{cm}^{-1}$  [43], which results in a bit lower D value than the experimental ones. Taking  $Dq = 1100 \, \mathrm{cm}^{-1}$ , we obtain  $D_{PT}$  close to the experimental values. Also, it is worth to consider how the changes in Dq affect the ZFSP D. Fig. 2 shows a clear linear

**Table 6** Calculated the ZFS parameter  $D_{PT}$  (in  $10^{-4}$  cm<sup>-1</sup>) and CF parameters (in cm<sup>-1</sup>) for Mn<sup>2+</sup> ion at the Cd<sup>2+</sup> (upper part) and NH<sub>4</sub><sup>+</sup> (lower part) sites in CAS.

		$D_{PT}$	$B_{20}$	$B_{40}$	B <sub>43</sub>
Calc. <sup>a</sup>	*	80.12	-946.1	-4073.9	-6459.5
	**	139.0	<b>-1351.3</b>	-5525.7	-8794.3
For $Dq = 1100 \text{ cm}^{-1}$		161.62	<b>-1471.7</b>	-6020.7	-9582.0
Calc. <sup>b</sup>	*	25.3	-581.5	-1810.2	-2870.2
	**	40.6	-830.6	-2455.3	-3907.6
Calc. <sup>c</sup>	*	2231	-2795.6	-24,787.0	-39,303.0
	**	4141	-3992.9	-33,621.0	-53,509.0
Calc. <sup>a</sup>	*	-81.2	2579.4	-1545.5	-2137.9
	**	-102.7	3334.0	-3075.5	-4245.3
For $Dq = 1100 \text{ cm}^{-1}$		-113.69	3631.0	-3348.4	-4621.9
Calc.b	*	-46.8	1585.4	-686.7	-949.9
	**	-61.3	2049.3	-1366.6	-1886.3
Calc. <sup>c</sup>	*	-227	7621.7	-9403.7	-13,008.0
	**	0.00059	9851.6	-18,713.0	-25,830.0

The notes (a) to (c) and \*, \*\* have the same meaning as in Table 5.



**Fig. 2.** Variation of ZFS parameter  $D_{PT}$  at both Cd<sup>2+</sup> and NH<sub>4</sub>+sites versus the cubic crystal-field parameter Dq for Mn<sup>2+</sup> centers in CAS.

**Table 7** Comparison of the calculated  $D_{SPM}$  and  $D_{PT}$  with the experimental  $D_{\text{exp}}$  for  $\text{Mn}^{2+}$  ions at  $\text{Cd}^{2+}$  and  $\text{NH}_4^+$  sites in CAS; all parameters in  $(10^{-4} \, \text{cm}^{-1})$ .

D parameter	Cd <sup>2+</sup> sites	NH <sup>4+</sup> sites
D <sub>SPM</sub>	158.9	9.6
$D_{PT}$	161.62	-113.69
$D_{\rm exp}$ [7]	$161.83 \pm 0.67$	
$D_{\rm exp}$ [5]	$-157.46 \pm 0.13$	
D <sub>exp</sub> [6]	$-159.7\pm0.5$	

relationship between  $D_{PT}$  and Dq at both sites. For comparison, the calculated  $D_{SPM}$  and  $D_{PT}$  values for  $\mathrm{Mn^{2+}}$  at  $\mathrm{Cd^{2+}}$  and  $\mathrm{NH_4^+}$  sites are tabulated in Table 7 together with the experimental  $D_{\mathrm{exp}}$ .

## 5. Conclusion

Zero-field splitting (ZFS) parameters (ZFSPs)  $b_2^0$  (=D) and  $b_4^q$  (q = 0, 3) of Mn<sup>2+</sup> in CAS crystal have been investigated using superposition model (SPM) and crystallographic data at the substitutional Cd<sup>2+</sup> and NH<sup>4+</sup> sites. The SPM results on  $b_2^0$  (=D) were confirmed, and its sign was determined as positive utilizing the fourth-order perturbation formula on the basis of the dominant spin-orbit coupling mechanism. The ZFS parameters obtained for Mn<sup>2+</sup> ion at Cd<sup>2+</sup> site using both approaches are in good agreement with the experimental values. Hence, it is indicated theoretically

that  $Mn^{2+}$  ions substitute for the  $Cd^{2+}$  ions in CAS crystal in addition to experimental suggestions.

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